#### Collapse of the attractive Bose-Einstein condensates

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#### BEC with repulsive interactions

- Predicted by Einstein for non-interacting bosons (1924)
- First experimental observation of BEC in dilute gases of alkali atoms with repulsive interactions (Anderson et al, 1995; Davis et al, 1995)

#### Experimental setup: cooling and trapping neutral atoms

- 1 Laser cooling, RF cooling:  $(T \sim 1 \mu K)$
- 2 Magnetic or optical trapping in 3D: parabolic anisotropic potential

$$V_{trap}(
ho,z)=rac{m\omega_{||}z^2}{2}+rac{m\omega_{\perp}
ho^2}{2}$$



# Particle interactions in BEC

Interaction parameter - scattering length

$$a_s: \quad U_0 \equiv \int U_{12}(\vec{r}_{12}) d\vec{r}_{12} = rac{4\pi \hbar^2 a_s}{m}$$

- Repulsive interaction  $\rightarrow$  positive  $a_s$
- Attractive interaction  $\rightarrow$  negative  $a_s$

Gross-Pitaevsky equation (GP):  $i\hbar\psi = -\frac{\hbar^2}{2m}\Delta\psi + V_{trap}\psi + U_0|\psi|^2\psi$ Trapped condensate is stable for repulsive interatomic interactions!

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Yes! If the particles number  $N < N_{cr}(\omega, U_0)$ .  $N > N_{cr}(\omega, U_0)$  - collapse!

#### Experiments with $Li^7$ ( $a_s = -29.2a_b$ ):

- I C.A. Sackett et al, "Measurements of Collective Collapse in a Bose-Einstein Condensate with Attractive Interactions", Phys. Rev. Lett. 82, 876 (1999).
- 2 J M. Gerton et al "Direct observation of growth and collapse of a Bose–Einstein condensate with attractive interactions", Nature 408, 692-695 (2000)

## Controlled BEC collapse

Feschbach resonance:  $a_s = a_s(B)$ 



Switching *a<sub>s</sub>* from positive to negative! - Controlled collapse.

#### 'Bosenova'' experiments:

- (JILA, Colorado, 2001): E. A. Donley et al, "Dynamics of collapsing and exploding Bose–Einstein condensates", Nature 412, 295-299 (2001)
- 2 (Australia, 2011): P. A. Altin et al, "Bosenova and three-body loss in a Rb-85 Bose-Einstein condensate", arXiv:1108.2561v1

## Collapse dynamics

Main features of collapse:

Missing atoms + burst atoms + jets + remnant condensate



Lack of the reliable theoretical model

## Theoretical models

Quantum many-body problem:

$$i\hbar\partial_t\hat{\Psi} = \left(-rac{\hbar^2}{2m}\Delta + V_{trap} + U_0\hat{\Psi}^\dagger\hat{\Psi}
ight)\hat{\Psi}$$

- Mean field approximation  $(N_c \gg 1)$ :  $\hat{\Psi} = \psi_0 \rightarrow \text{GP}$  equation:  $i\hbar\partial_t\psi_0 = \left(-\frac{\hbar^2}{2m}\Delta + V_{trap} + U_0|\psi_0|^2\right)\psi$
- At  $r \sim a_s$  GP breakes down!

#### Corrections to the GP equation:

- 3 body recombination (dimer + 1 escaping atom)  $i\hbar\partial_t\psi_0 = \left(-\frac{\hbar^2}{2}\Delta + V_{trap} + U_0|\psi_0|^2 + iK_3|\psi_0|^4\right)\psi_0$
- **Beyond mean-field:** fluctuations  $\hat{\Psi} = \psi_0 + \hat{\chi}$

$$i\hbar\partial_t\hat{\chi} = \left(-\frac{\hbar^2}{2m}\Delta + V_{trap}\right)\hat{\chi} + U_0\left(2|\psi_0|^2\hat{\chi} + \psi_0^2\hat{\chi}^\dagger\right)$$

Condensate/uncondensate particles density  $n_{cond} = |\psi_0|^2$ ,  $n_{unc}(\vec{r}) = \langle \chi^{\dagger}(\vec{r})\chi(\vec{r}) \rangle$ 

#### Energy estimates

What effect is more important? 3-body loss  $\propto n_{cond} a_s^3 \ll 1$  gaseous parameter.

Our idea: the uncondensate particles generation is crucial for  $r \gg a_s!$ The particles energies are too high! Kinetic energy:  $T_{kin} \sim \frac{\hbar^2}{m|a_s|}$  Critical

temperature in trap: 
$$T_0^{trap} = \hbar \bar{\omega} N^{\frac{1}{3}}, \qquad \bar{\omega} = \sqrt[3]{\omega_{||}\omega_{\perp}^2}$$

Experiment	Atoms	Part. Number	Cond. Temp. T <sub>0</sub> <sup>harm</sup> K	Kin. Energy, K
Donley et. al., 2001 Nature	Rb <sup>85</sup>	6E+3	1E-9	2E-5
Altin, Dennis et al, 2011 ar	Rb <sup>85</sup>	4E+4	8E-9	4E-5

Uncondensate particles generation ⇔ Collapse stability Numerical experiment is nedeed. (S. Wuster, J. J. Hope, and C. M. Savage, PRA 2005)

# Elliptic deformations instability

Strong quasiclassical collapse in NLSE (Zakharov, Kuznetsov; JETP 1986)

- Generalization elliptic collapse  $|\psi_0|^2(\rho, z, t) = \lambda^2 \left[ 1 - \left(\frac{\rho}{a(t)}\right)^2 - \left(\frac{z}{b(t)}\right)^2 \right].$
- Scaling parameters: effective potential  $V_{eff}(a,b) = -\frac{\lambda^2}{a^2b}$ .
- The collapse with a(t) = b(t) is unstable!



## Summary

- The problem of the theoretical description of the attractive BEC collapse is not solved after decade.
- The decoherence effects are underestimated in the existing theoretical approaches.
- Elliptic type instability in strong NLSE collapse is found. Jets (?).